

WHICH FUTURE FOR NEUTRINOLESS DBD ?



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AN ENDLESS RESEARCH FIELD

- How much does a neutrino weigh ?
- What is the mass ordering (hierarchy)
- Is neutrino a Majorana or Dirac particle
- Do more (sterile) neutrinos exist ?
- Do neutrinos violate CP ?
- Can we observe the CNB (a picture of a universe 1 second old)

MAJORANA VS. DIRAC



$$V_L^M \xleftarrow[\text{Lorentz}]{\text{CPT}} V_R^M$$

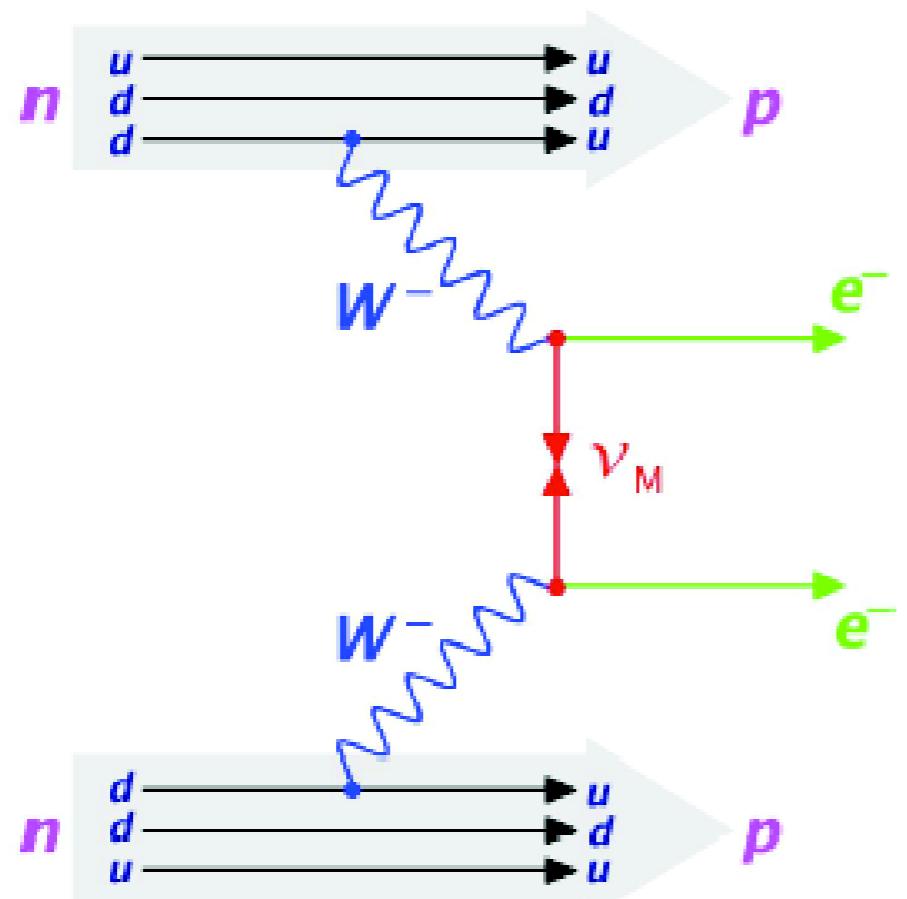
Majorana



$$\begin{array}{ccccc} & & V_L^D & \xleftarrow{\text{Lorentz}} & V_R^D \\ & \uparrow & & & \downarrow \\ & \text{CPT} & & & \text{CPT} \\ & \downarrow & & & \uparrow \\ & \overline{V}_R^D & \xleftarrow{\text{Lorentz}} & \overline{V}_L^D & \end{array}$$

Dirac

NEUTRINOLESS DOUBLE BETA DECAY



$$(T_{1/2}^{0\nu})^{-1} = G^{0\nu}(Q, Z) |M^{0\nu}|^2 \left(\frac{\langle m_{\beta\beta} \rangle}{m_e} \right)^2$$

Phase Space Factor Nuclear Matrix Element

Only if:

Majorana Neutrinos
Massive Neutrinos

If observed:

Proof of the Majorana
nature of Neutrino
Indication of mass scale

how long should we wait to see it (if any) ?

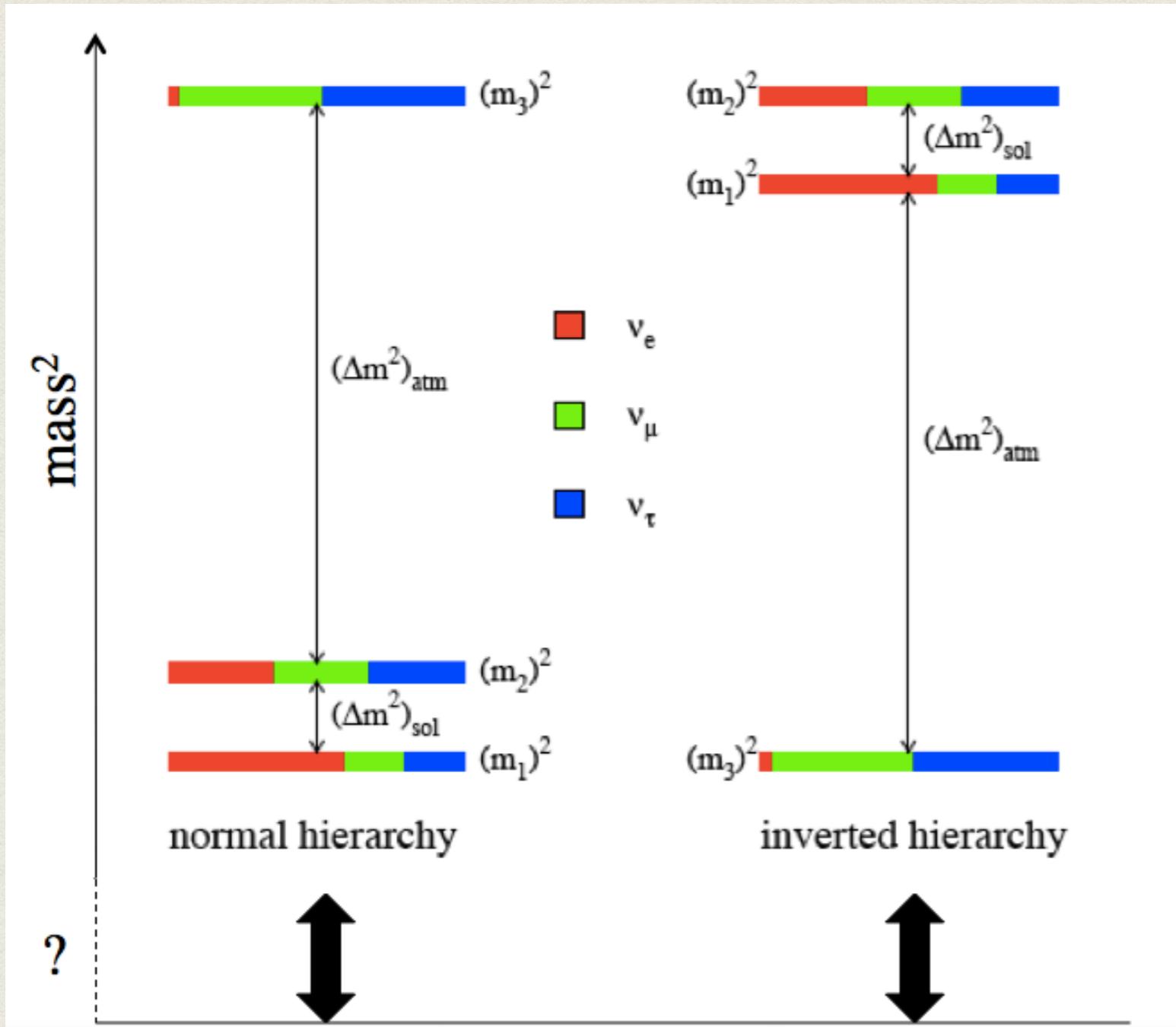
parameter containing
the **physics**

$$1/\tau = G(Q, Z) |M_{\text{nucl}}|^2 \langle M_{\beta\beta} \rangle^2$$

what the **experimentalists**
try to measure

what the **nuclear theorists**
try to calculate

MASS HIERARCHY

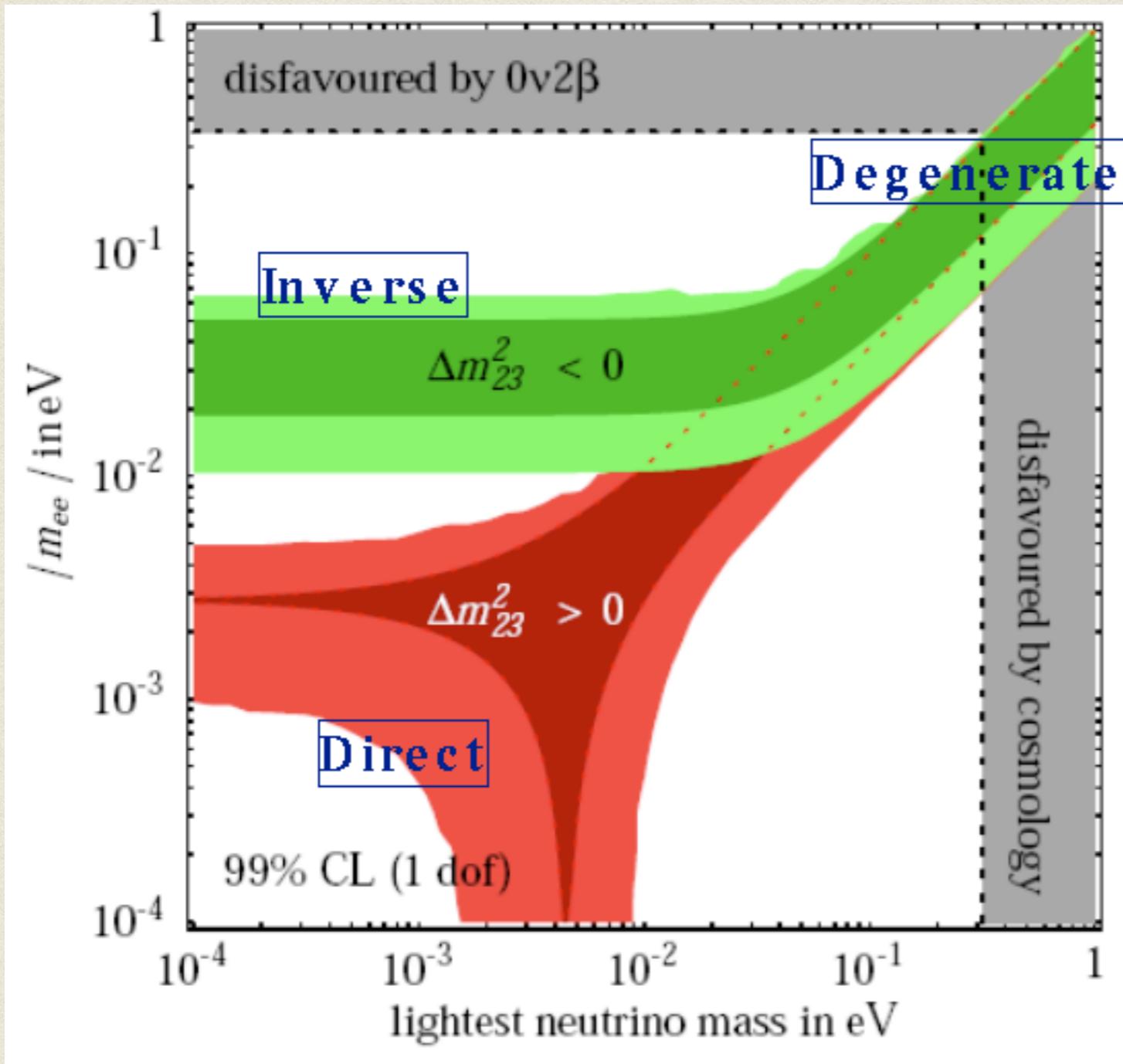


1-2 ordered by matter

3 is free

degeneracy belong to an other type of experiments

NEUTRINOLESS DBD INVERTED OR DIRECT ?



Three region

- 1) Degenerate : few 100meV
- 2) Inverse: few 10meV
- 3) Direct: few meV

just on the back
of the envelope

$$\left[T_{1/2}^{0\nu} \right]^1 = C \cdot \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$$

$C \sim 10^{-12} \text{ y}^{-1}$, $m_e \sim 500 \text{ keV}$, $m_{\beta\beta} \sim 20 \text{ meV}$

$$\tau_{1/2}^{0\nu} > 10^{25} \text{ y}$$

[universe life $15 \cdot 10^9 \text{ y}$, Avogadro number $6 \cdot 10^{23}$]

The name of the game

expected
number of
 $\beta\beta 0\nu$ events

$$S = \frac{M \cdot N_A \cdot a}{W} \cdot \ln(2) \cdot \frac{t}{T_{1/2}^{0\nu}} \cdot \varepsilon$$

detector mass isotopic abundance live time efficiency
molecular mass / /
 ββ0ν half-life

mean number of
background counts
around the Q-value

$$B = b \cdot M \cdot \Delta E \cdot t$$

background rate in
counts/keV/kg/y energy resolution
(detector FWHM)
 / /
 detector mass live time

Sensitivity is S/\sqrt{B}

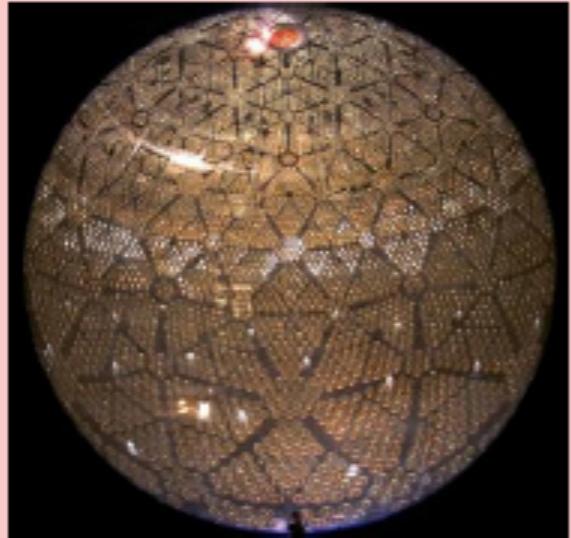
Sensitivity $\propto K \sqrt{\frac{M \cdot t}{B \cdot \Delta E}}$ (i.a. • ϵ)

$$m_{\beta\beta} \propto \sqrt{(1/\tau)}$$

To gain a factor 10 you need 10000 !

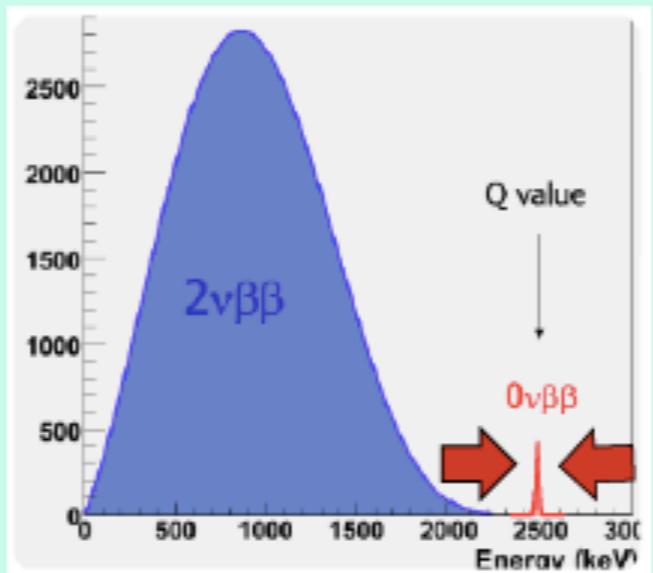
meaning :

The “Brute Force” Approach



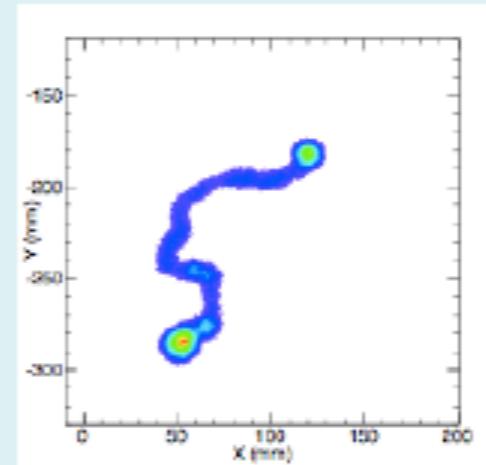
focus on the numerator
with a **huge amount**
of material
(often sacrificing
resolution)

The “Peak-Squeezer” Approach



focus on the denominator
by **squeezing down ΔE**
(various technologies)

The “Final-State Judgement” Approach

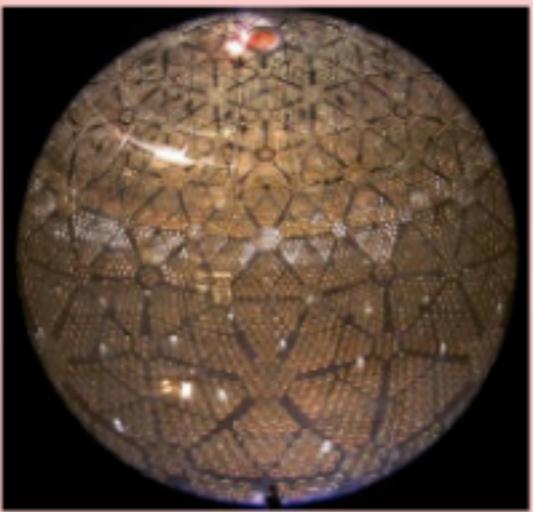


try to make the
background zero by
tracking or
tagging

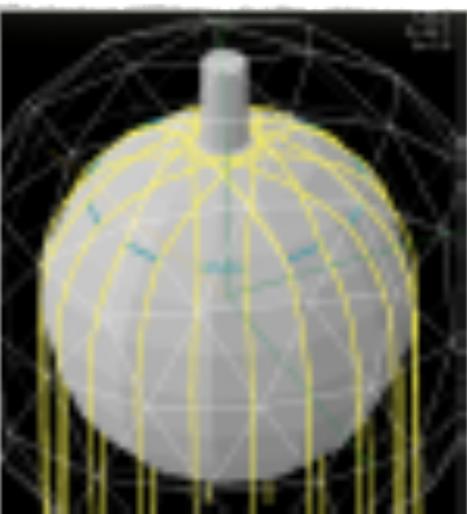
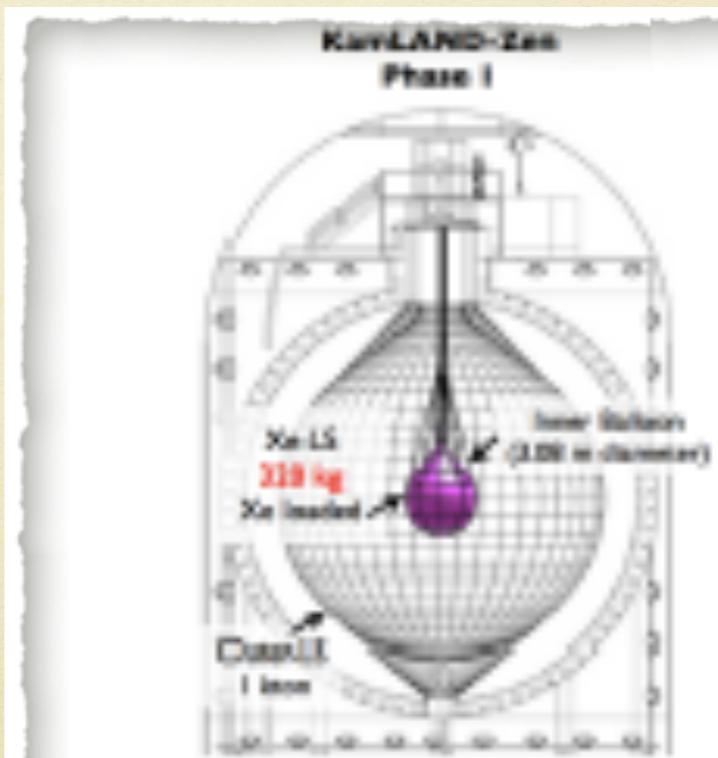
or better make the right cocktail of all of the above

the state of the art: brute force

The “Brute Force” Approach



focus on the numerator
with a **huge amount**
of material
(often sacrificing
resolution)



SNO+
 (^{130}Te)

KamLAND-Zen
 (^{136}Xe)

Kamland-Zen

Search for Majorana Neutrinos near the Inverted Mass Hierarchy Region with KamLAND-Zen

$$T_{1/2}^{0\nu} > 1.07 \times 10^{26} \text{ yr (90\% C.L.)}$$

535 days livetime

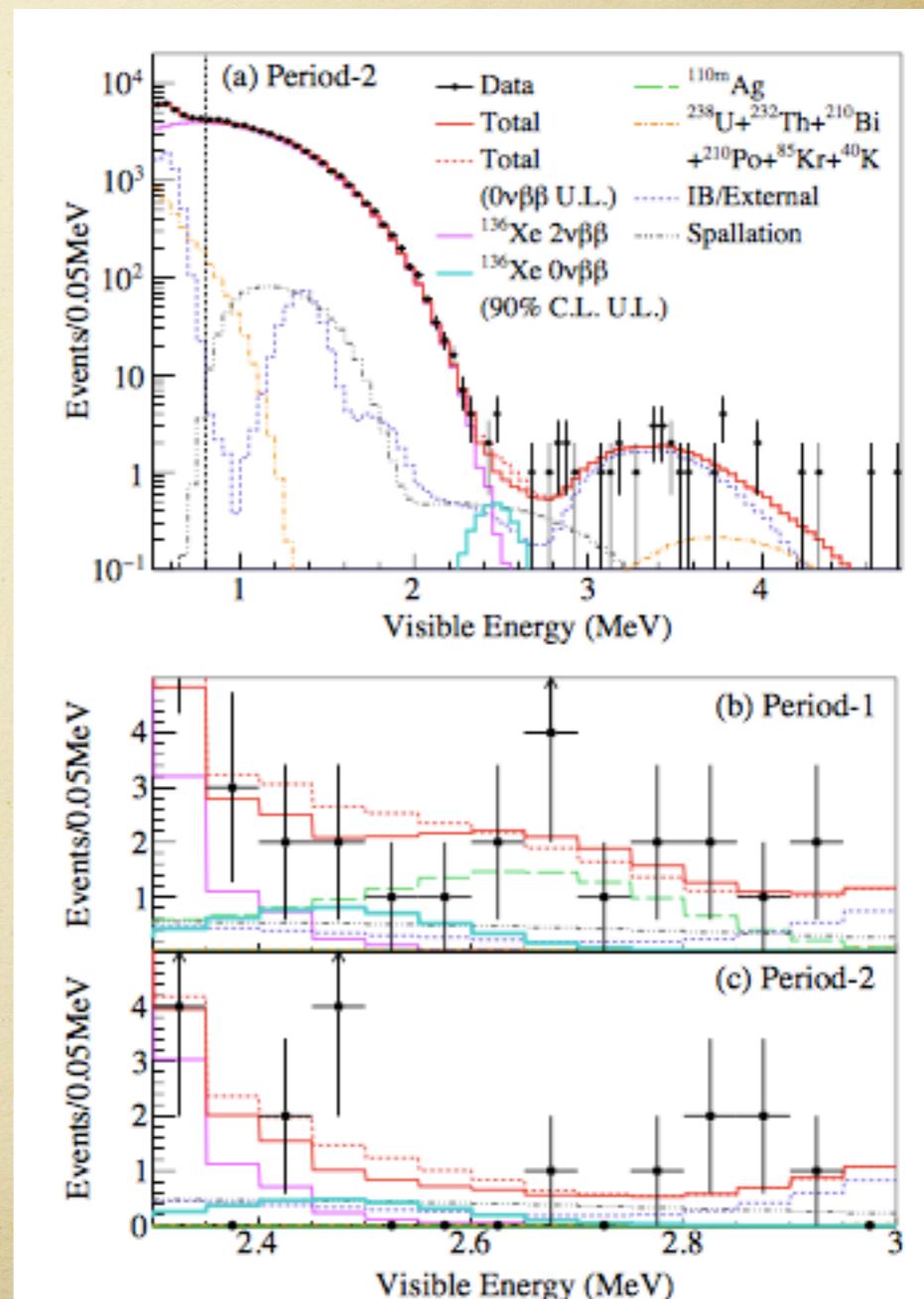
3700 mole*year

504 kg*year

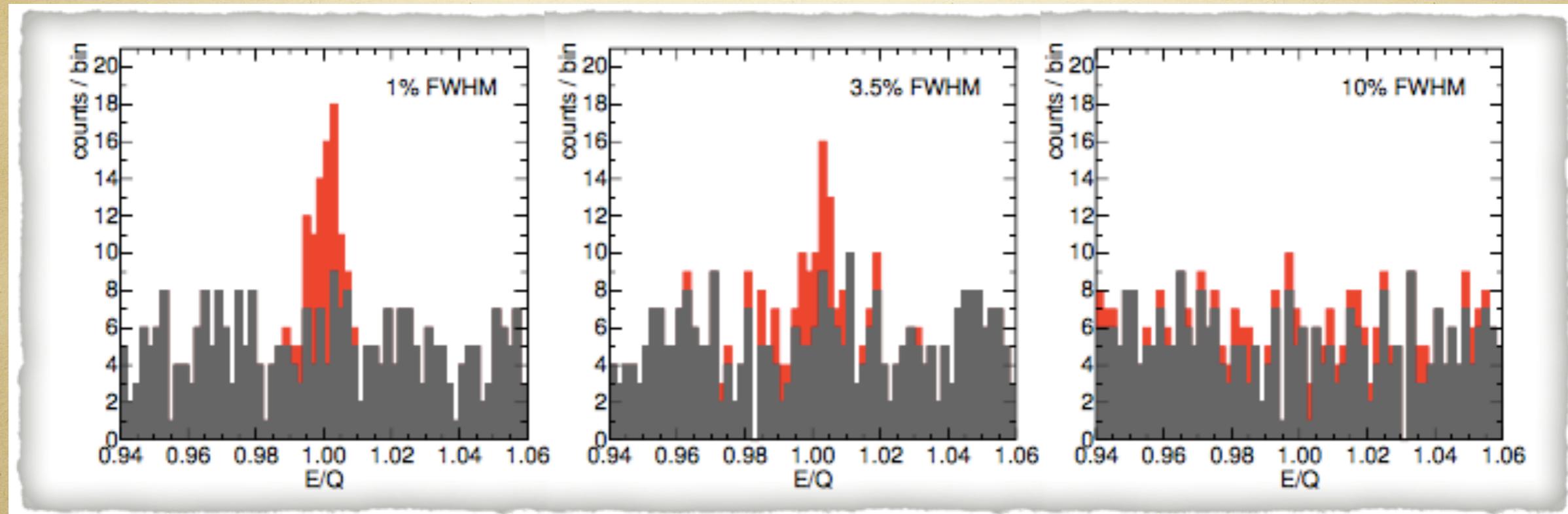
From the limit on the ^{136}Xe decay half-life, we obtain a 90% C.L. upper limit of $m < (61 - 165) \text{ meV}$ assuming the axial coupling constant

$$g_A \approx 1.27$$

	Period-1 (270.7 days)		Period-2 (263.8 days)	
Observed events	22		11	
Background	Estimated	Best-fit	Estimated	Best-fit
$^{136}\text{Xe } 2\nu\beta\beta$	-	5.48	-	5.29
Residual radioactivity in Xe-LS				
^{214}Bi (^{238}U series)	0.23 ± 0.04	0.25	0.028 ± 0.005	0.03
^{208}Tl (^{232}Th series)	-	0.001	-	0.001
^{110m}Ag	-	8.5	-	0.0
External (Radioactivity in IB)				
^{214}Bi (^{238}U series)	-	2.56	-	2.45
^{208}Tl (^{232}Th series)	-	0.02	-	0.03
^{110m}Ag	-	0.003	-	0.002
Spallation products				
^{10}C	2.7 ± 0.7	3.3	2.6 ± 0.7	2.8
^6He	0.07 ± 0.18	0.08	0.07 ± 0.18	0.08
^{12}B	0.15 ± 0.04	0.16	0.14 ± 0.04	0.15
^{137}Xe	0.5 ± 0.2	0.5	0.5 ± 0.2	0.4



effect of energy resolution

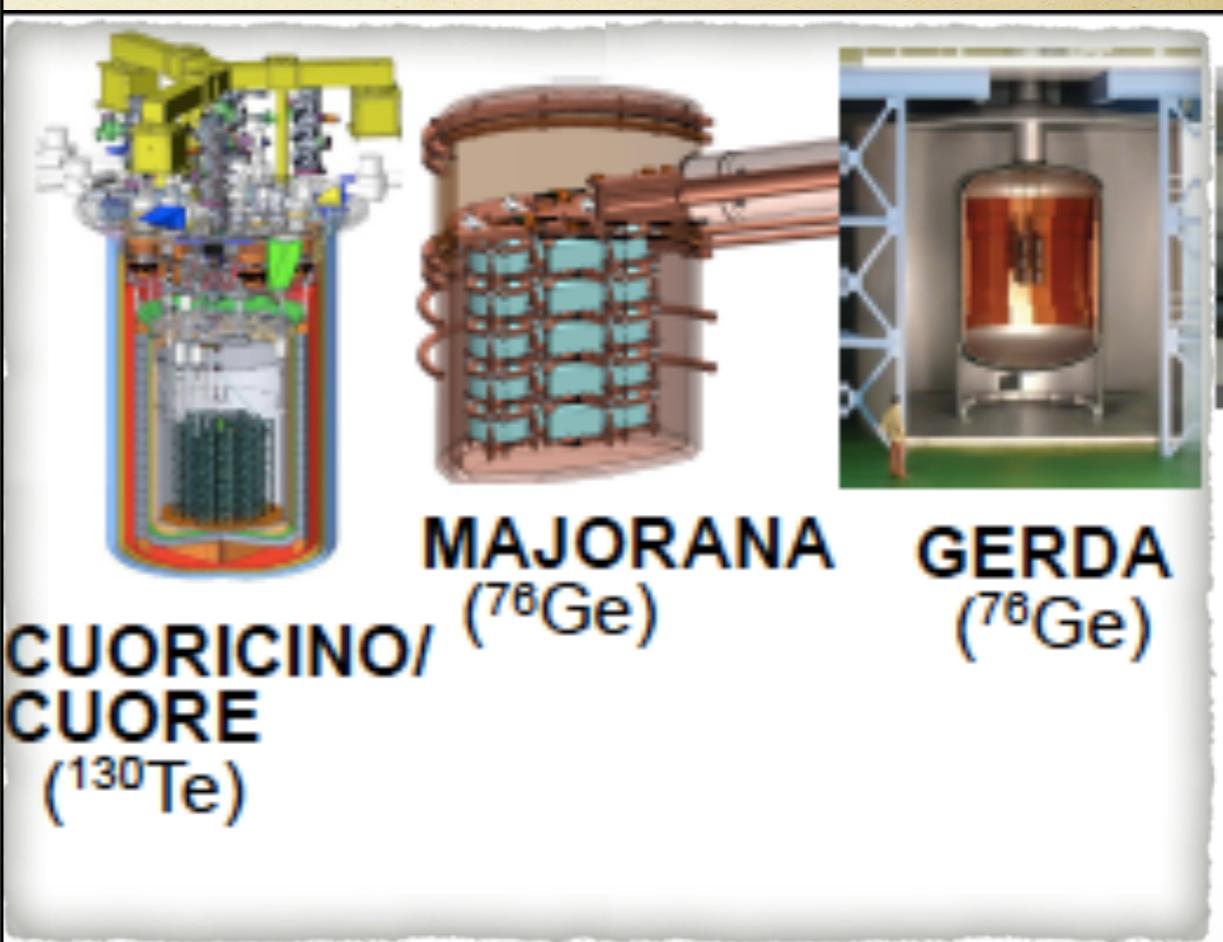
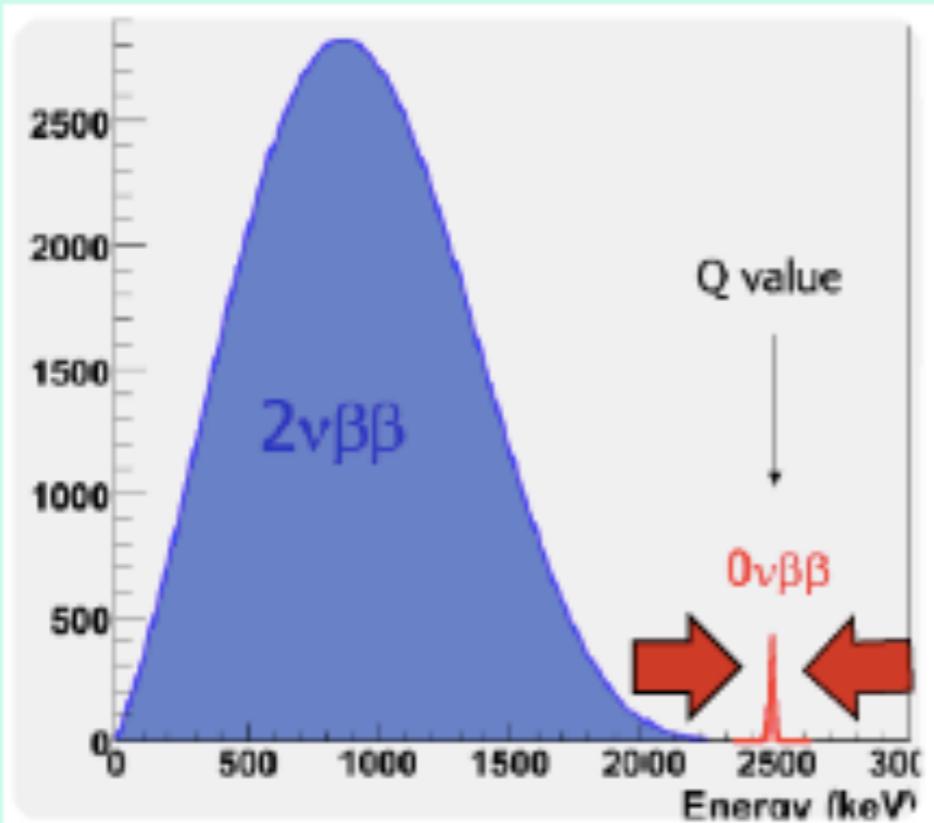


$S = 50$ events

$B = 1$ count/keV

the state of the art: peak squeezer

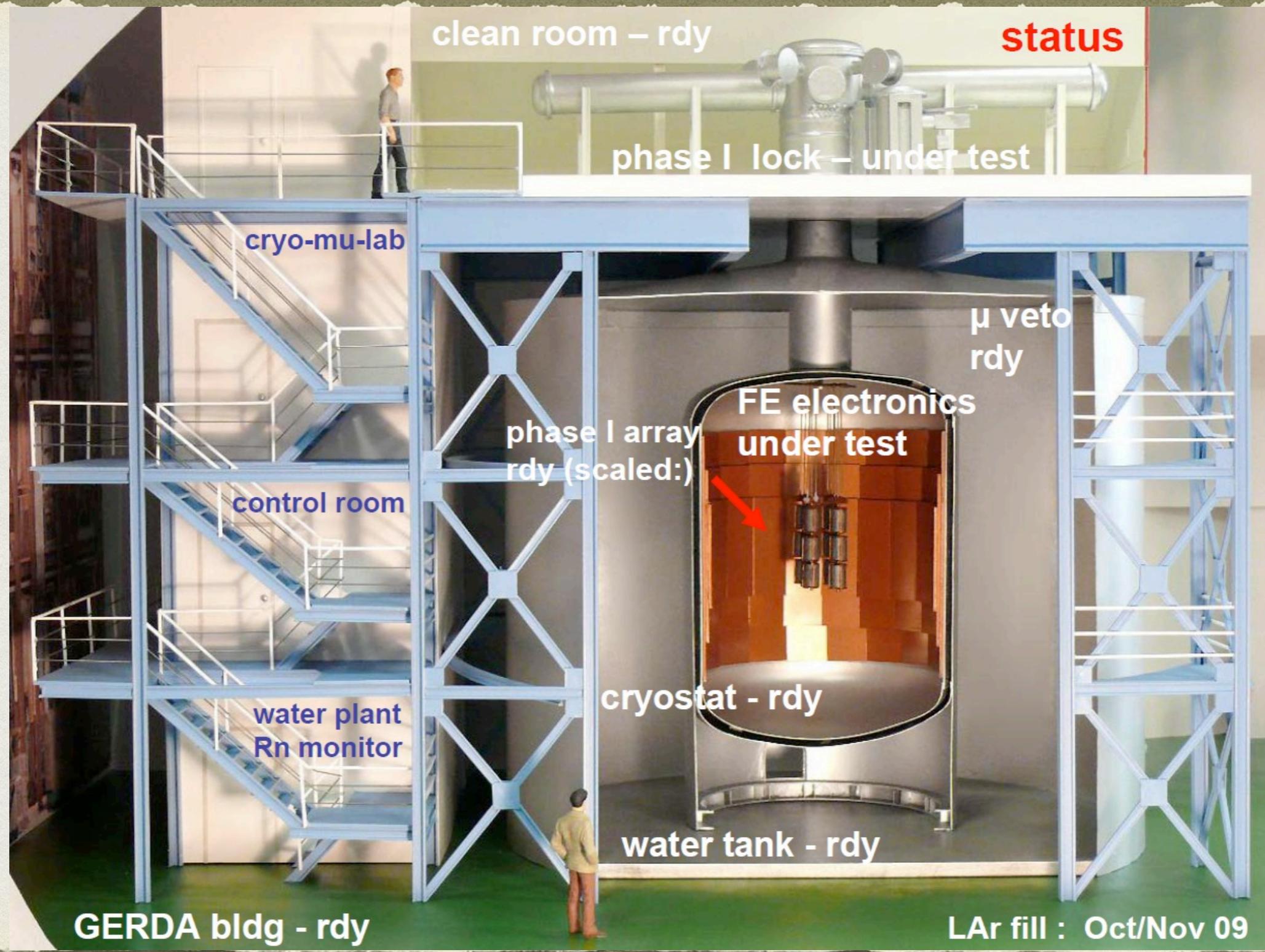
The “Peak-Squeezer” Approach



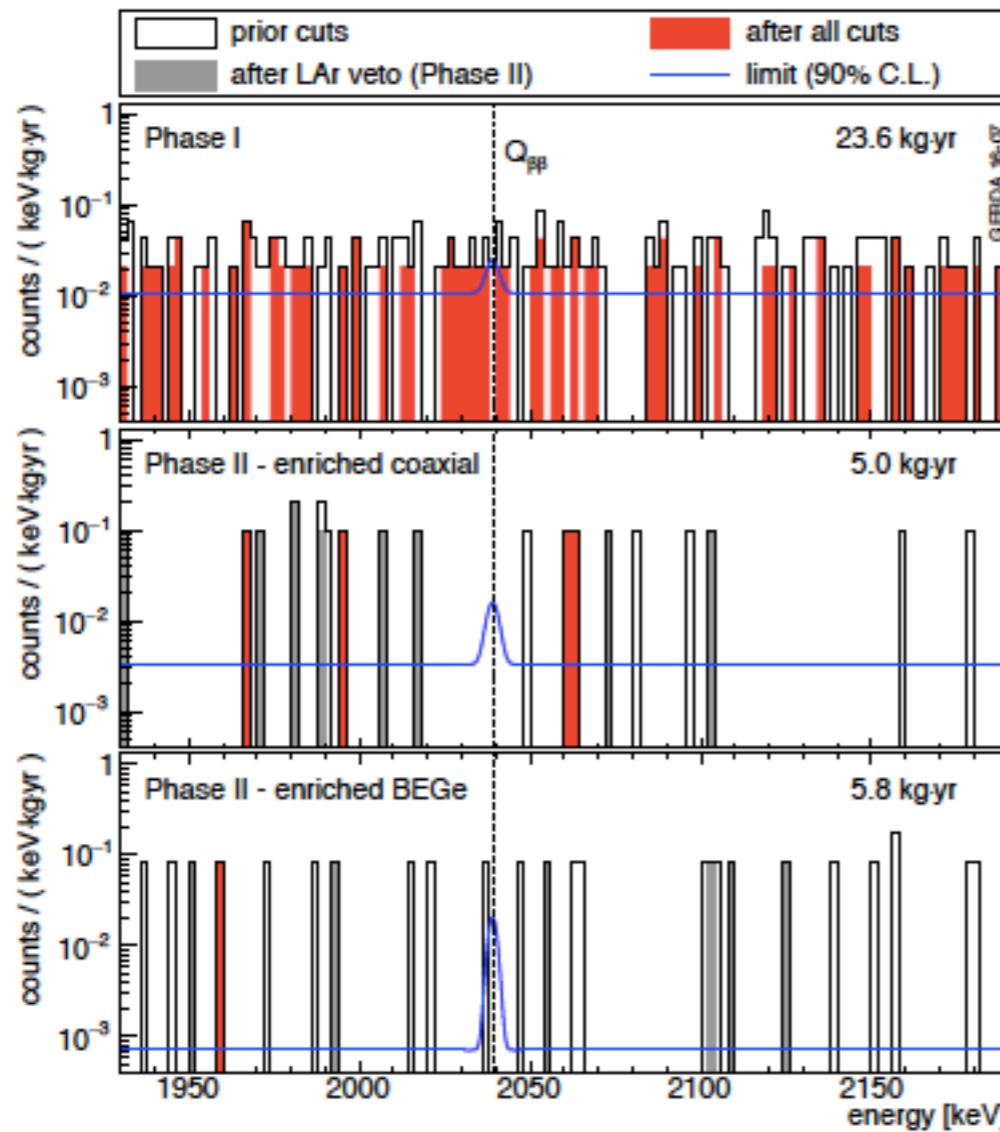
LNGS LEADS THE HUNT

- GERDA with a spectacular limit
- CUORE taking data
- CUPID paving the way for a new generation aiming at $B=0$

GERDA



THE BEST RESULT IN TOWN



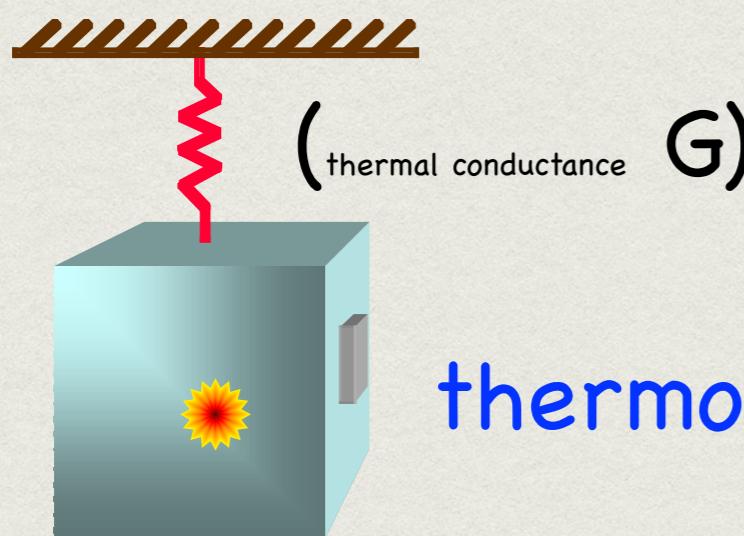
$$T_{1/2}^{0\nu} > 5.3 \cdot 10^{25} \text{ yr.}$$

We expect only a fraction of a background event in the energy region of interest (1 FWHM) at design exposure of 100 kg·yr. GERDA is hence the first **background free** experiment in the field. Our sensitivity grows therefore almost linearly with time instead of by square root like for

the GERDA half-life sensitivity of $4.0 \cdot 10^{25}$ yr for an exposure of $343 \text{ mol}\cdot\text{yr}$ is similar to the one of Kamland-Zen for ^{136}Xe of $5.6 \cdot 10^{25}$ yr based on a more than 10-fold exposure of $3700 \text{ mol}\cdot\text{yr}$ [9].

CUORE PRINCIPLE

heat sink



(T_0) Basic Physics: $\Delta T = E/C$

$$C(T) = \beta \frac{m}{M} \left(\frac{T}{\Theta_D} \right)^3$$

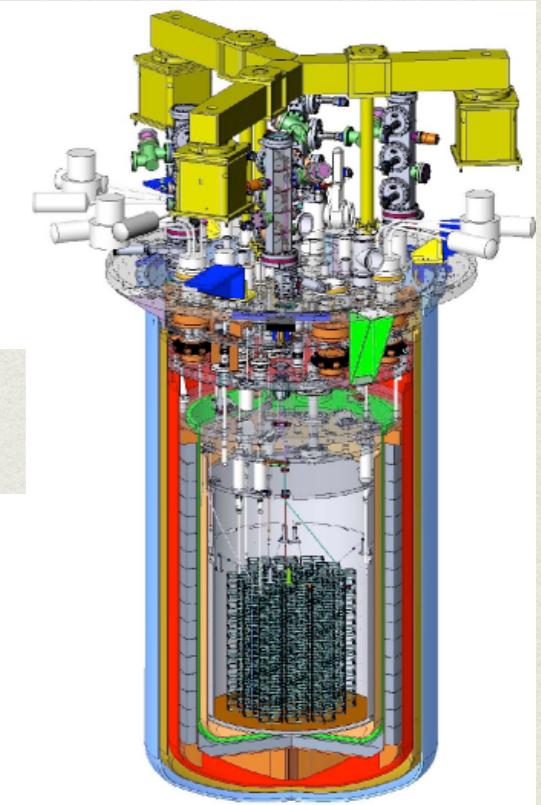
$\beta\beta$ atom x-tal

$$\Delta T(t) = \frac{\Delta E}{C} \exp\left(-\frac{t}{\tau}\right)$$

$T_0 \sim 10$ mK

$C \sim 2$ nJ/K ~ 1 MeV/0.1 mK

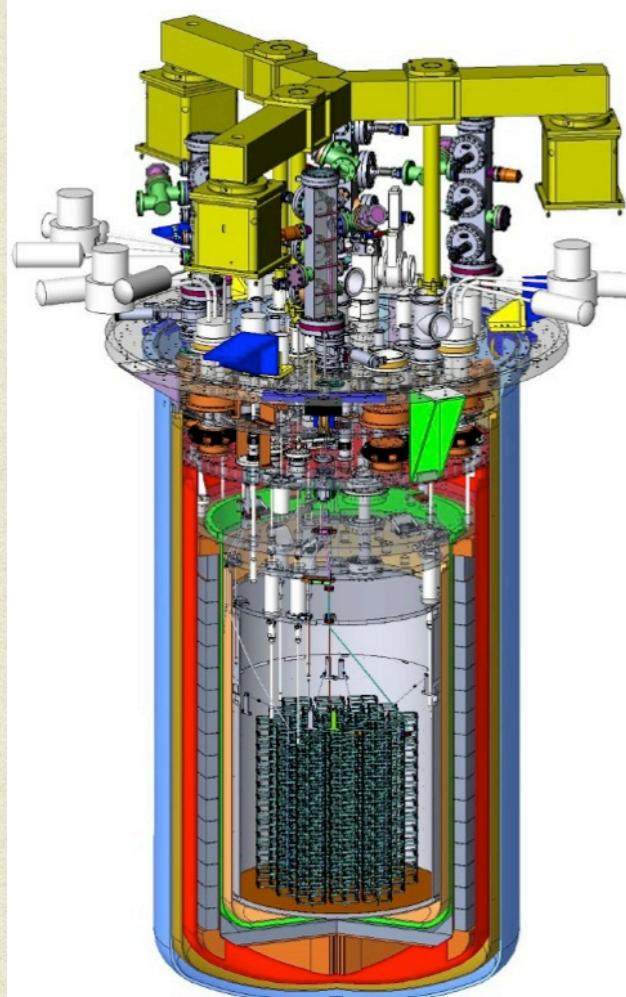
$G \sim 4$ pW/mK



CUORE

Cryogenic Underground Observatory for Rare Events

Searching for neutrinoless double beta
decay of ^{130}Te



200 Kg

^{130}Te

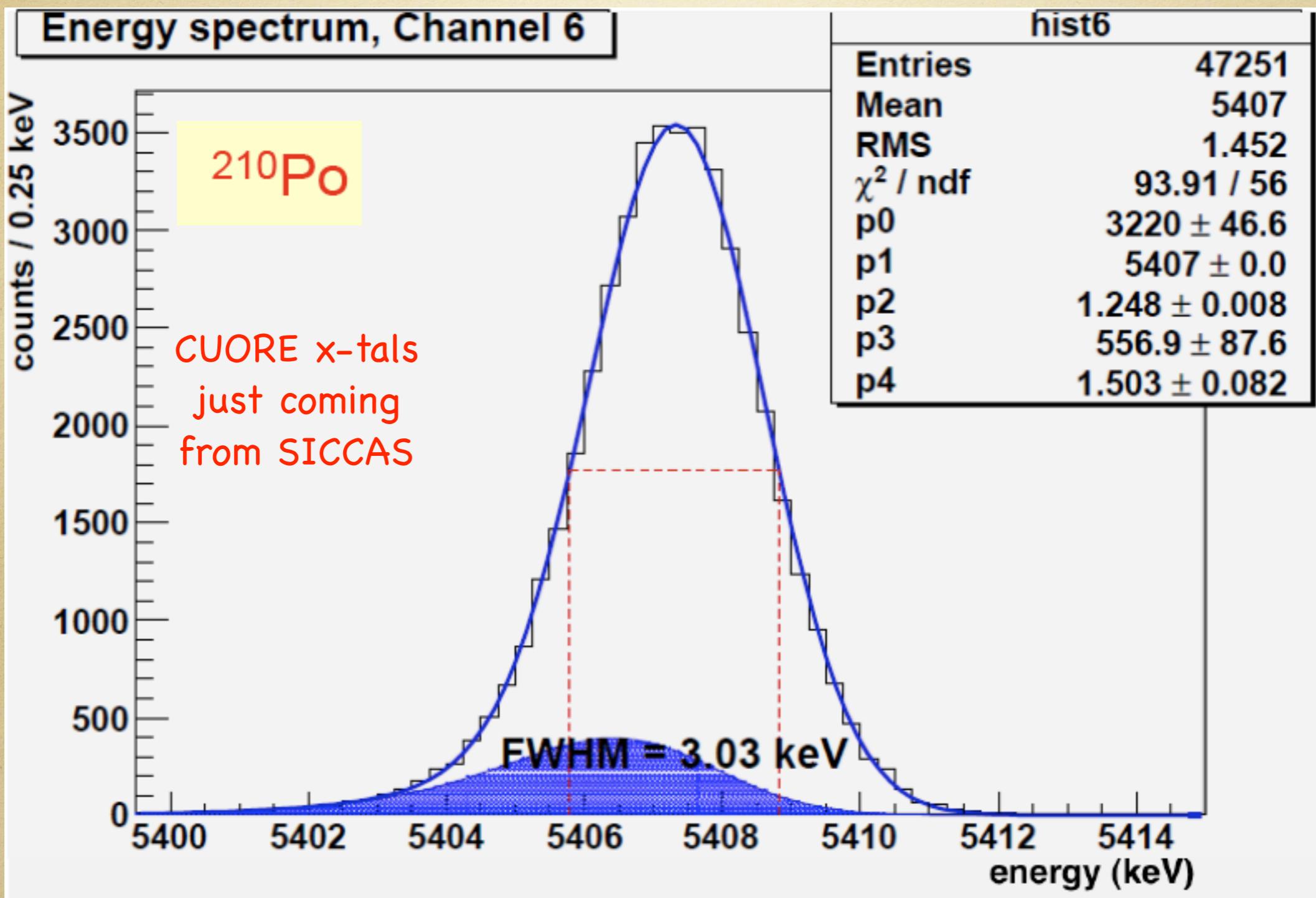


CUORE Hut

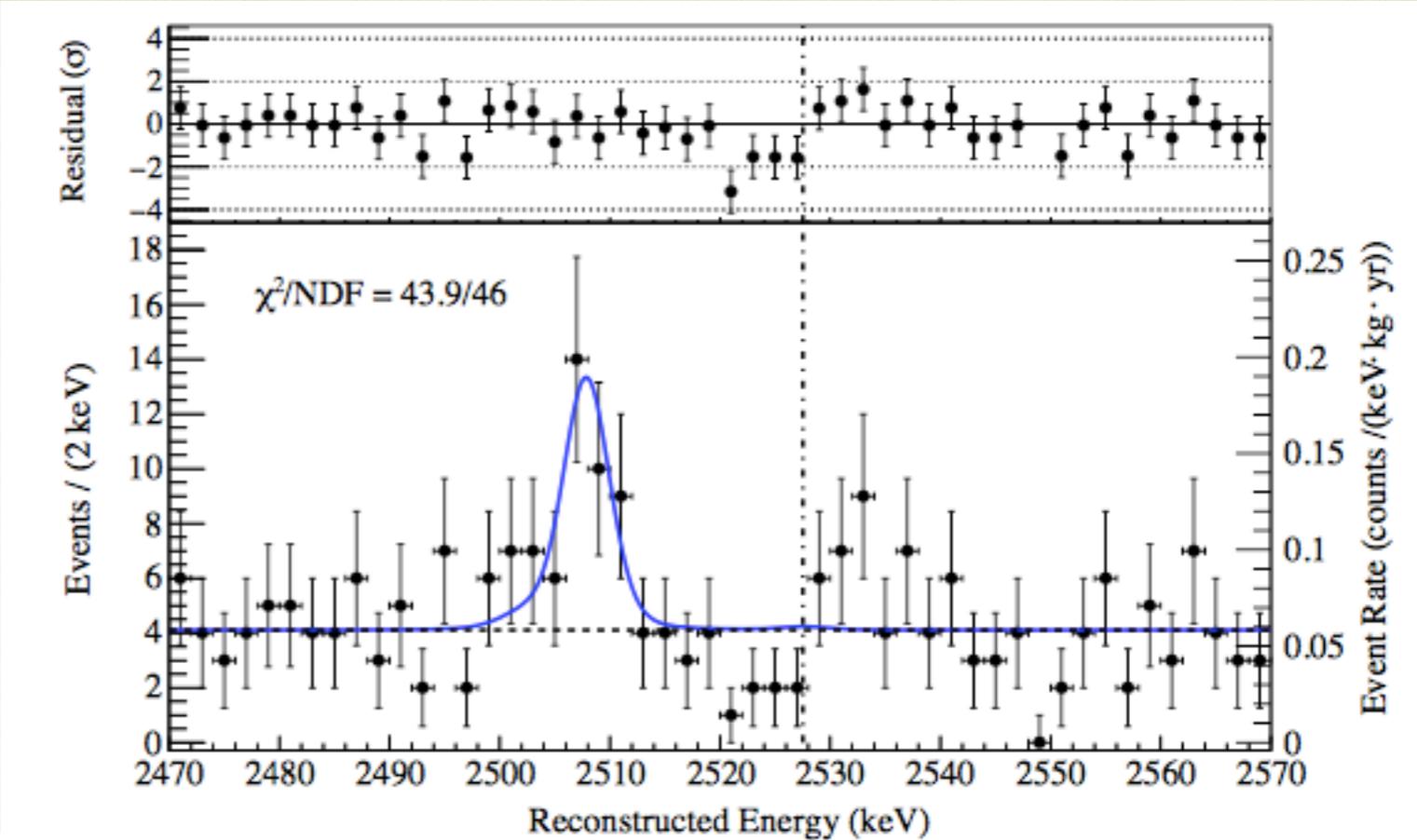
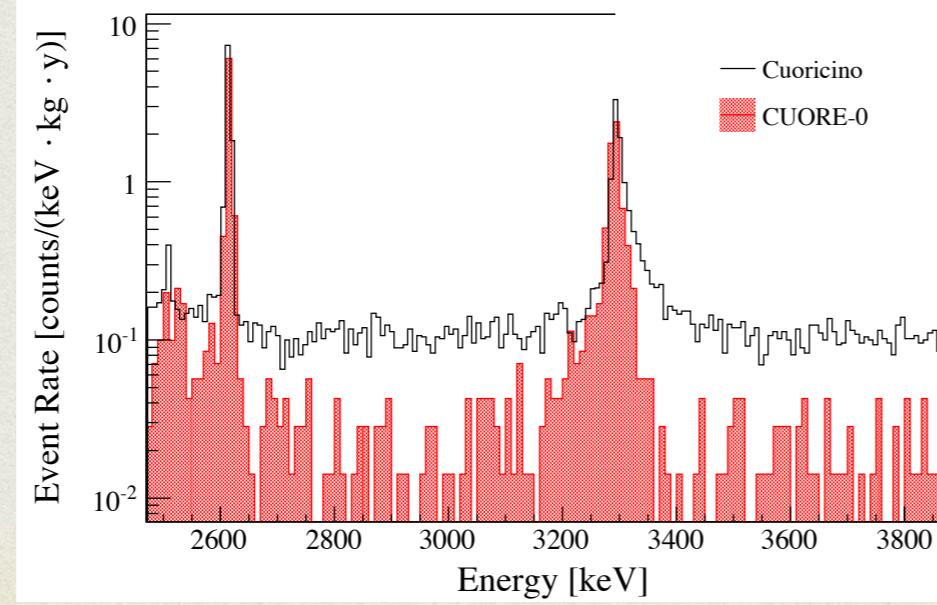


Expected 5 Years sensitivity:
 $T_{1/2} = 2.1 \times 10^{26} \text{ y}$, $m_{\beta\beta} = 41\text{--}95 \text{ meV}$
background counting rate
 $10^{-2} \text{ c/keV/kg/y}$

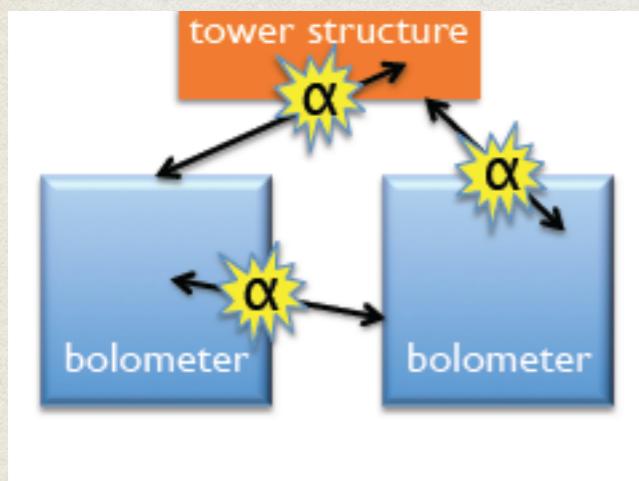
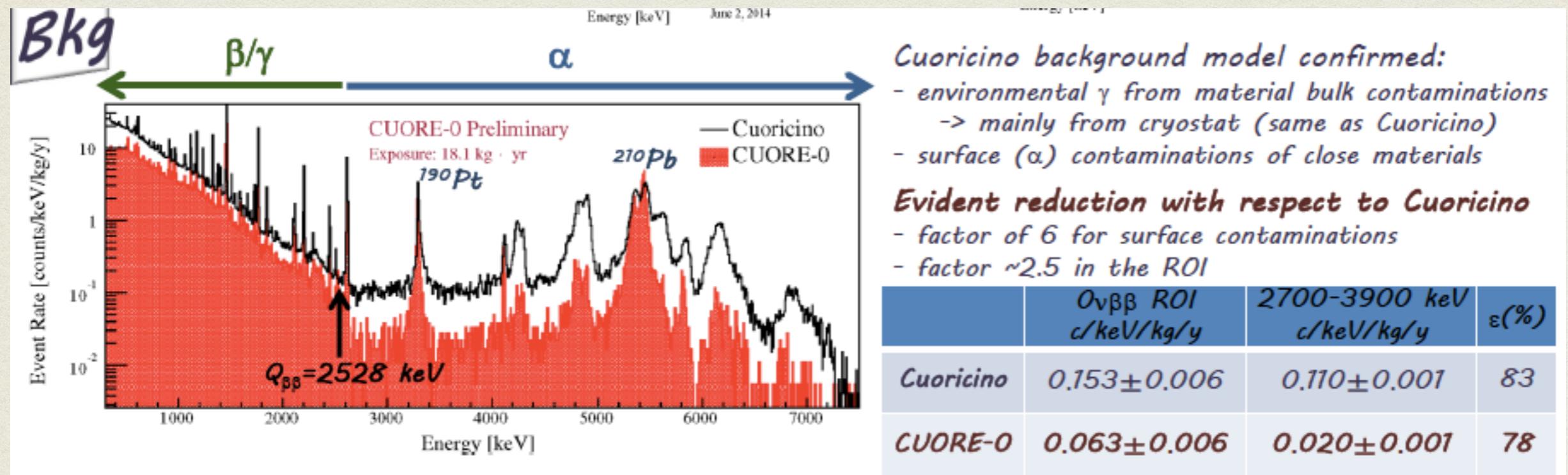
how much can you squeeze ?



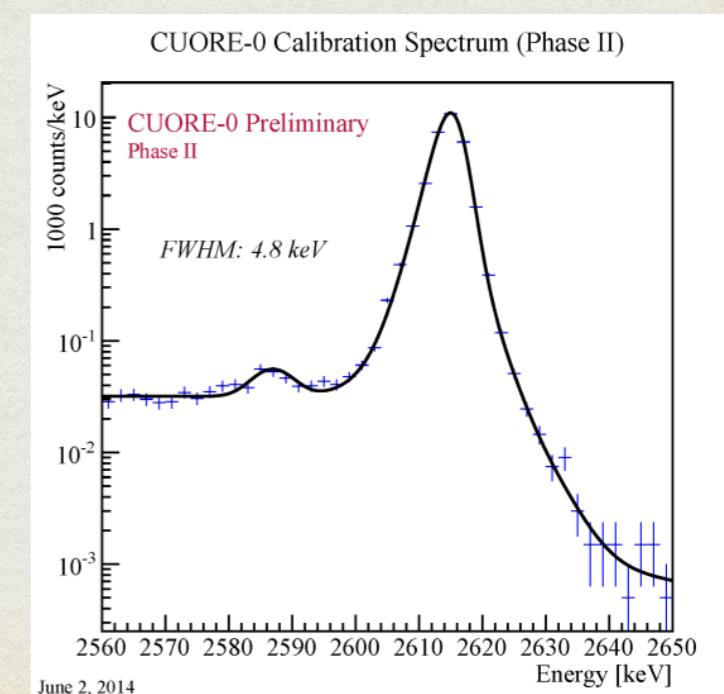
EXPERIENCE FROM CUORE-O



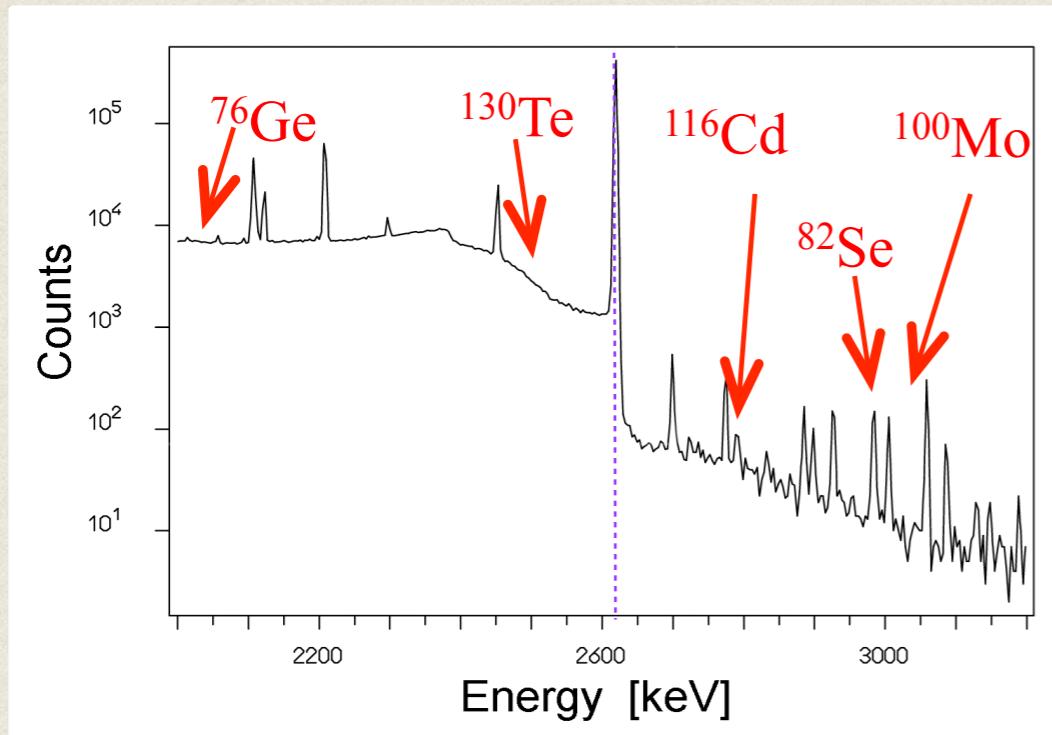
BUT WILL STOP AT



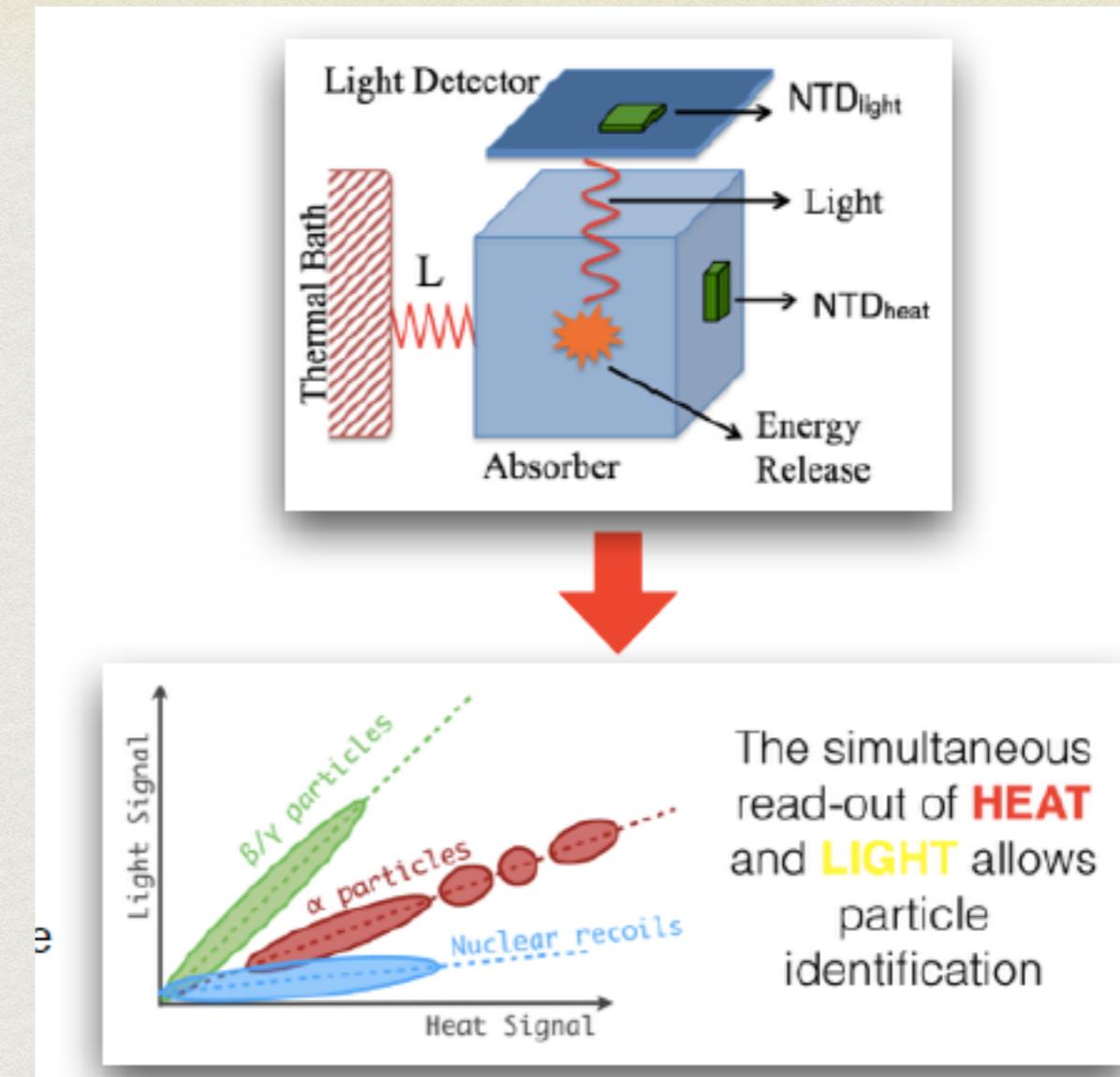
irreducible if you do not tell alphas from gammas in spite of formidable energy resolution



AN OPTION FOR A BRIGHT FUTURE

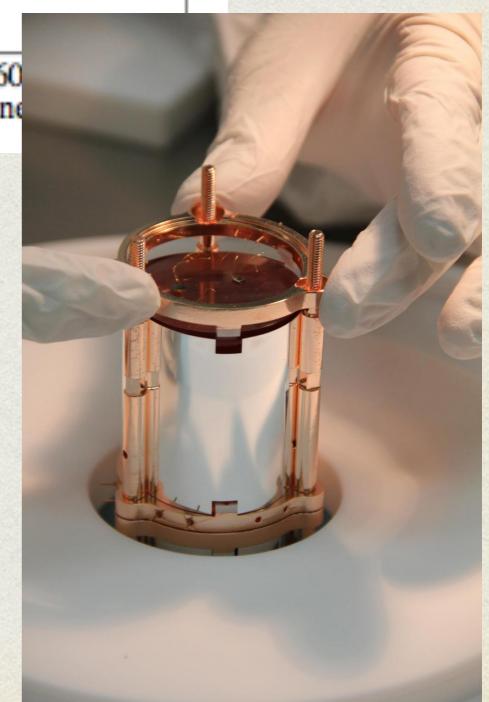
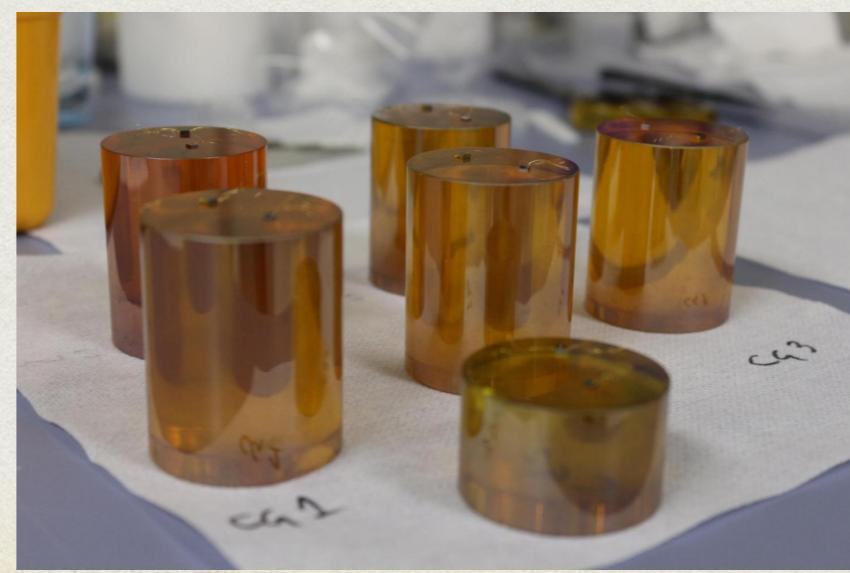
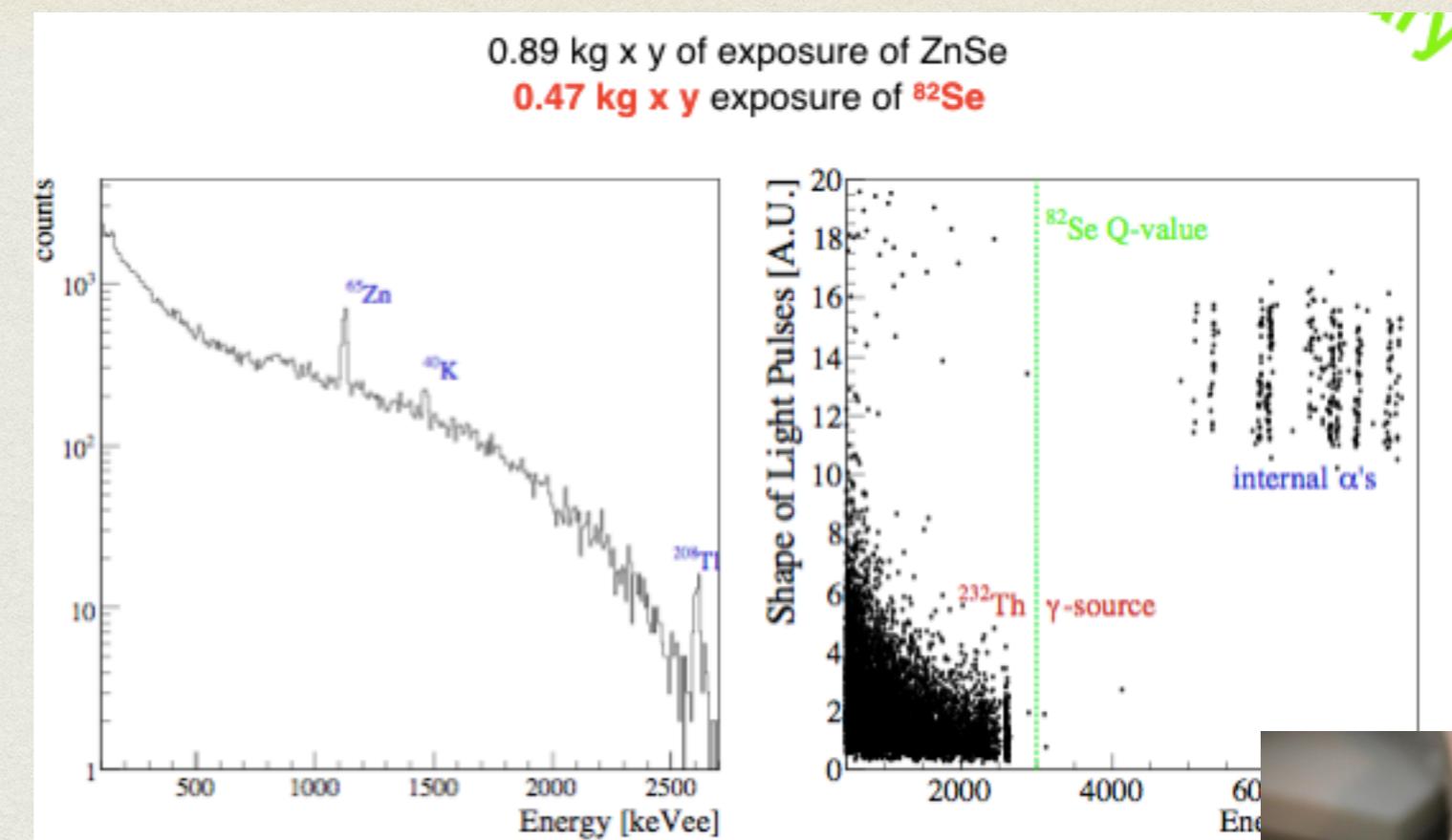
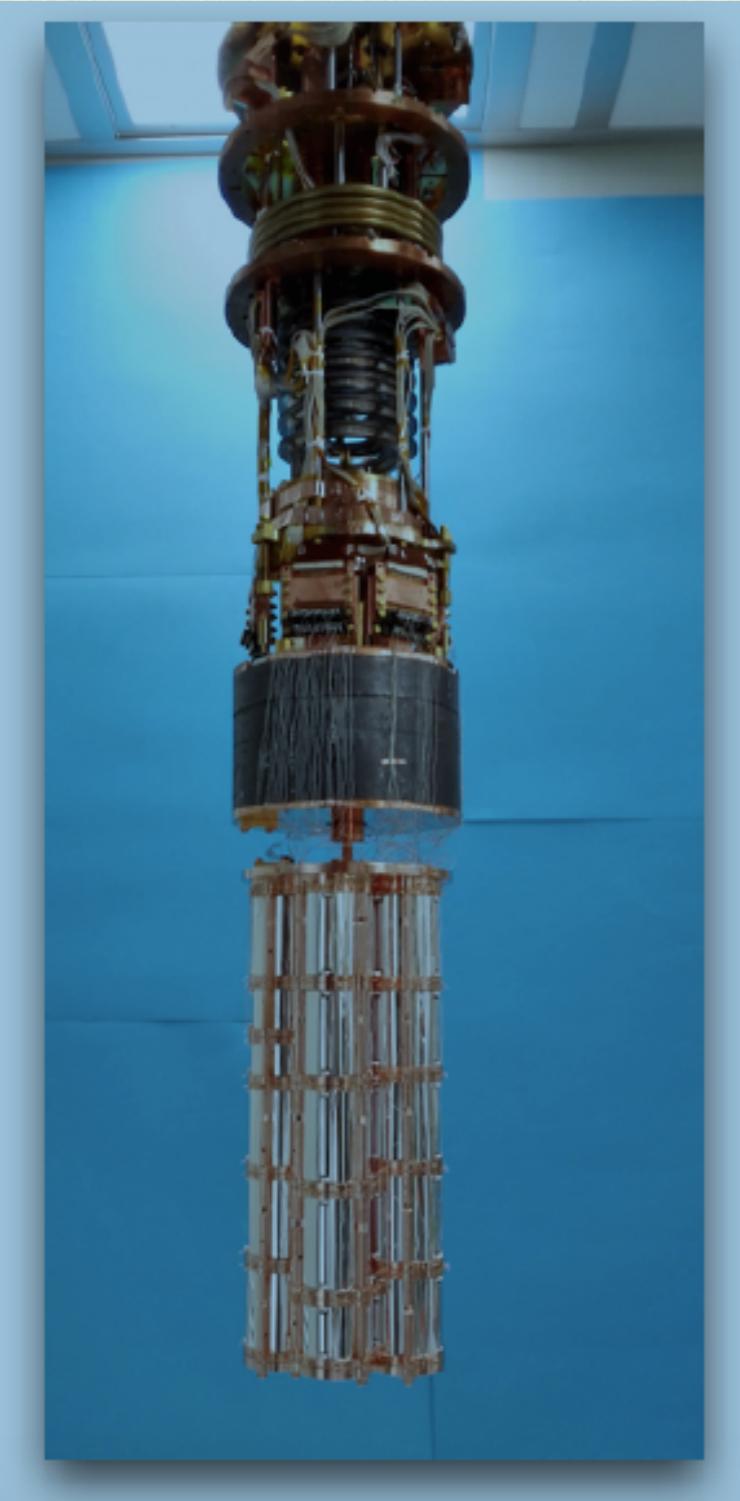


A **background-free experiment** is possible:
α-background: identification and rejection
β背景: $\beta\beta$ isotope with large Q-value

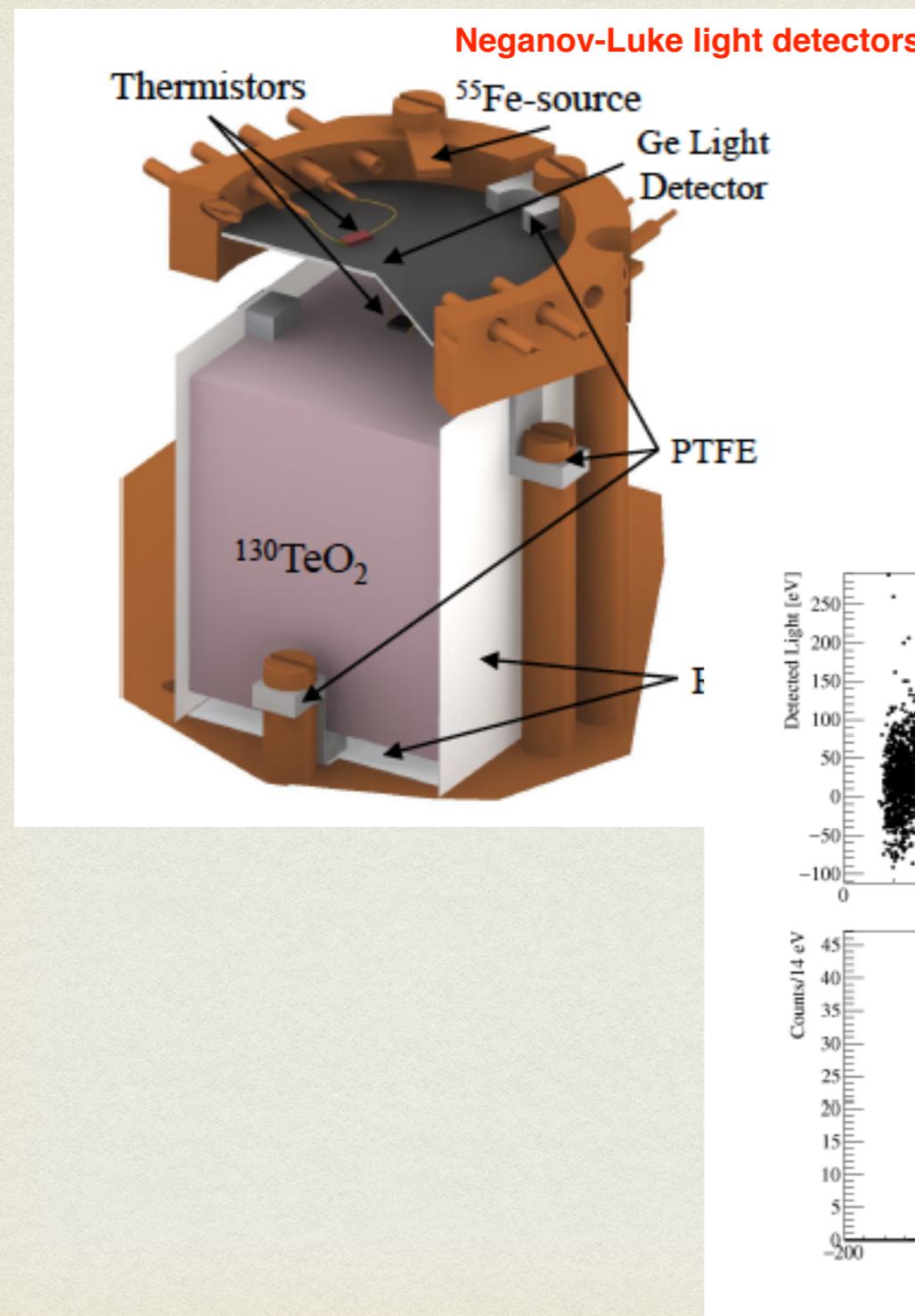


The simultaneous read-out of **HEAT** and **LIGHT** allows particle identification

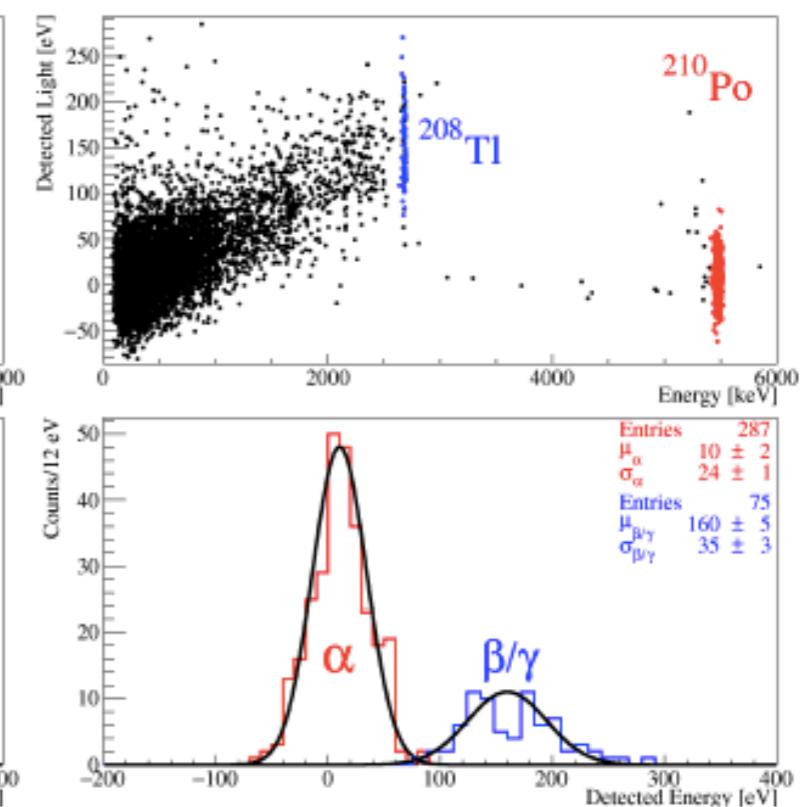
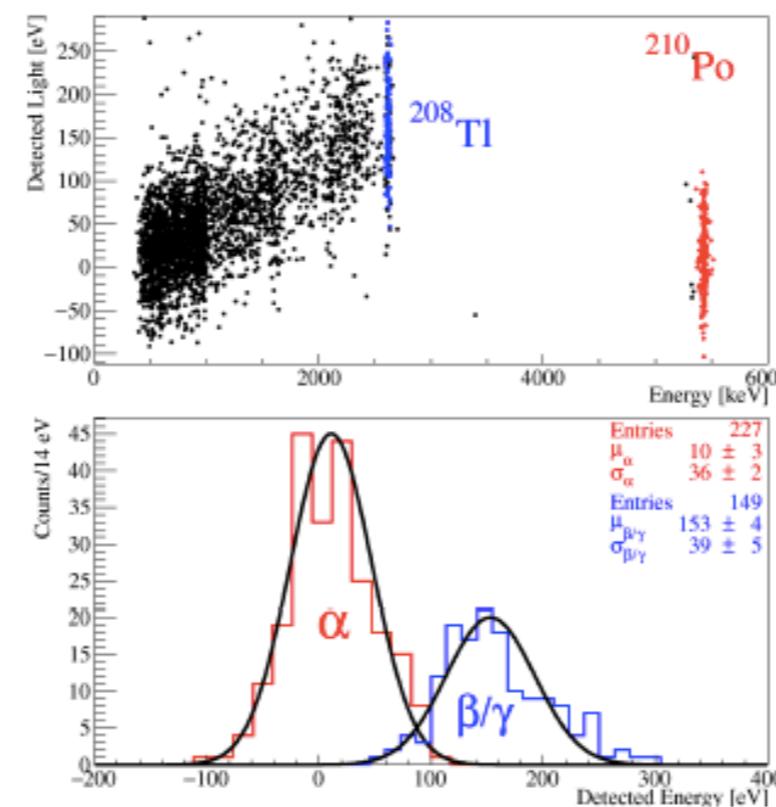
CUPID-o (LUCIFER)



YET ANOTHER OPTION

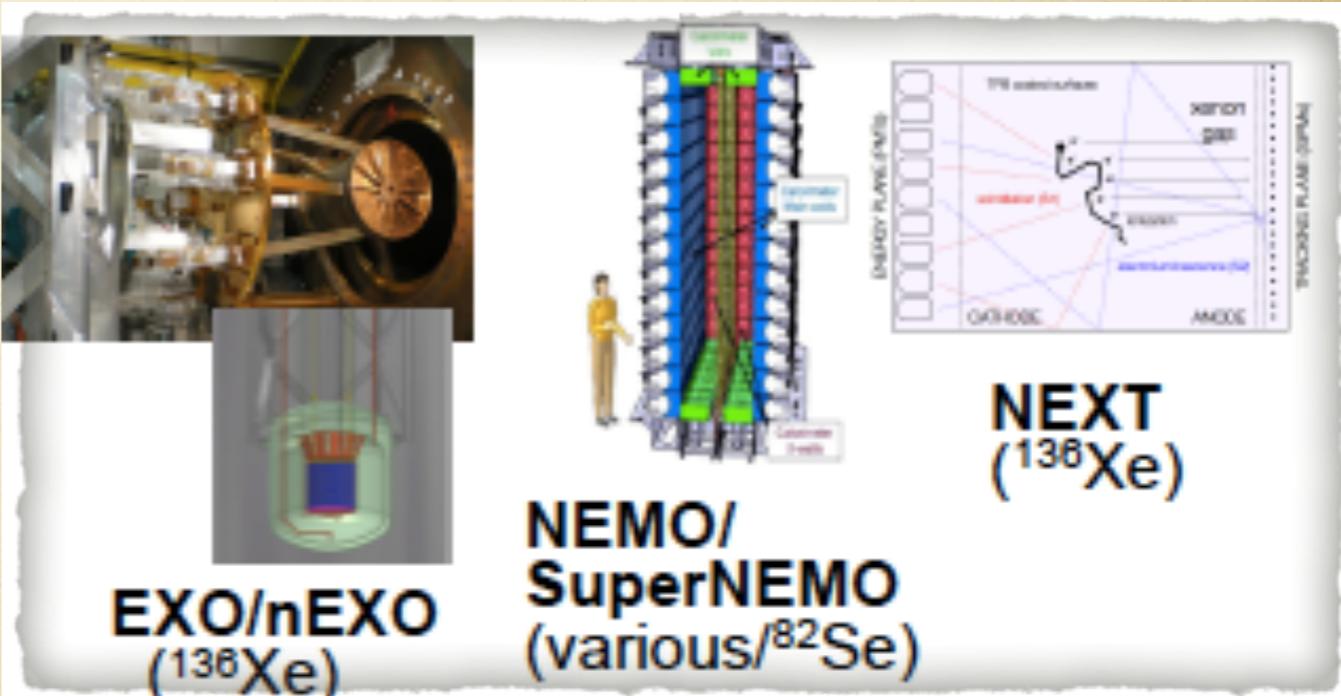
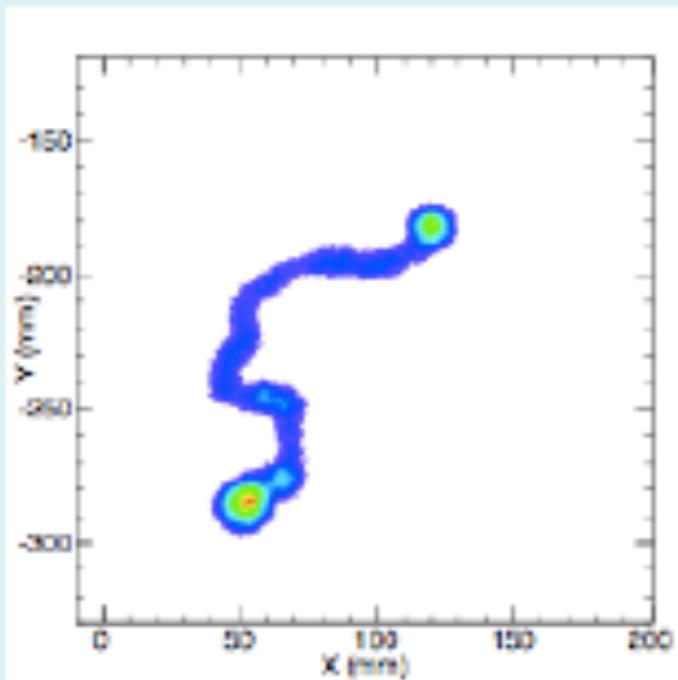


TeO₂ does not scintillate
Use Cherenkov effect

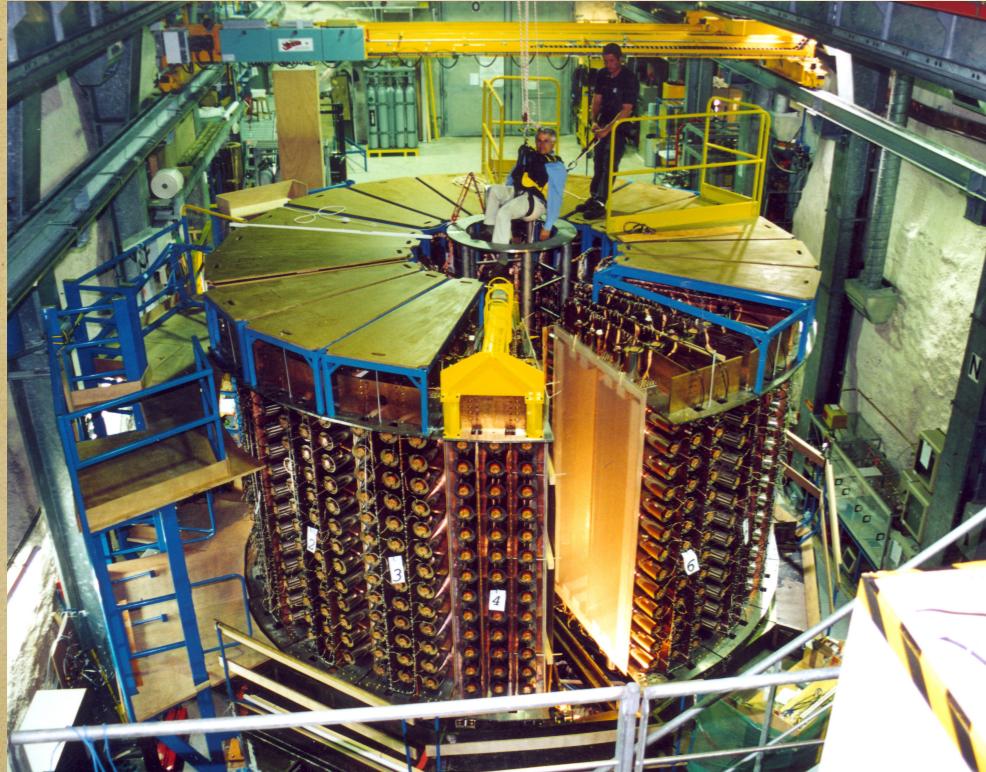


the state of the art: tracking

The “Final-State Judgement” Approach



nicely working but...



Source: 10 kg of $\beta\beta$ isotopes
cylindrical, $S = 20 \text{ m}^2$, $e \sim 60 \text{ mg/cm}^2$

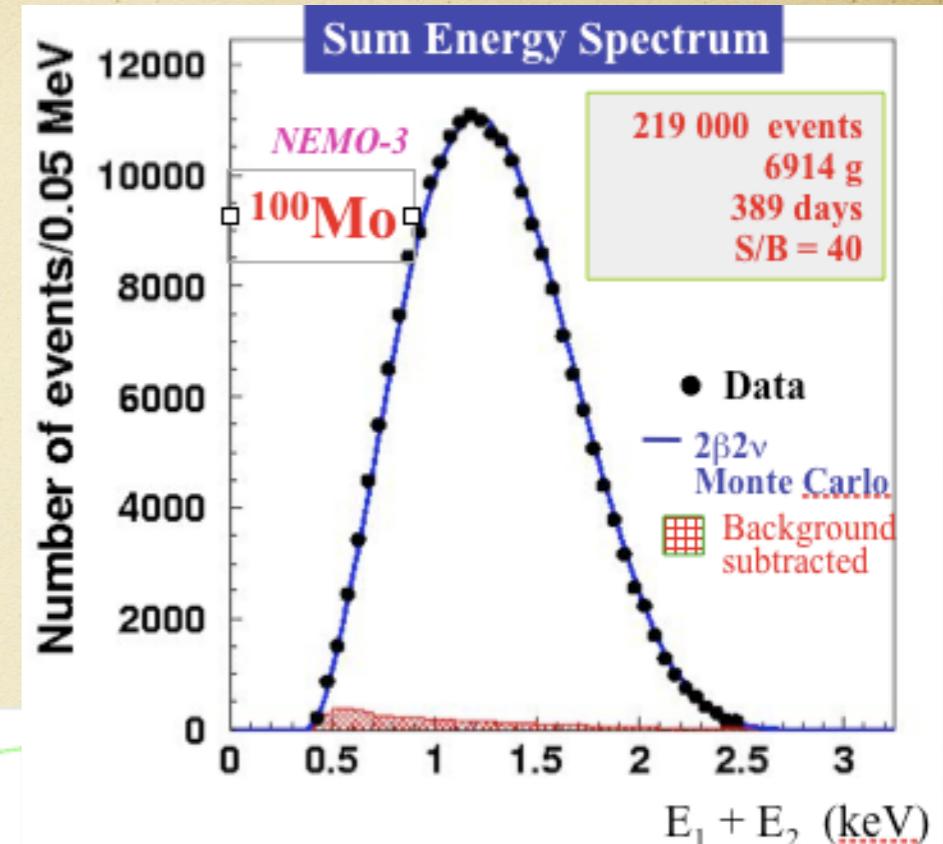
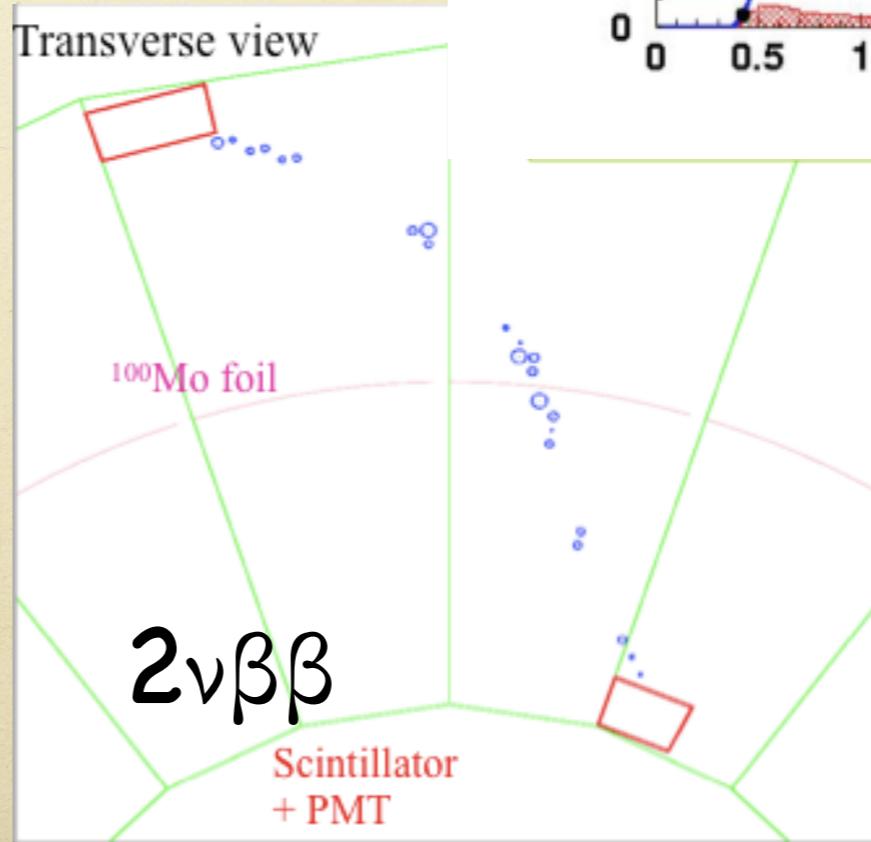
Tracking detector:

drift wire chamber operating
in Geiger mode (6180 cells)

Gas: He + 4% ethyl alcohol + 1% Ar + 0.1% H_2O

Calorimeter:

1940 plastic scintillators
coupled to low radioactivity PMTs



which way ?

- ▷ increase isotopic abundance close to 1 (linear though costly)
- ▷ increase M a lot (square root, say 1 ton)
- ▷ decrease B (get zero background and get rid of the square root !)
- ▷ get an extraordinary good energy resolution (remember we are talking of a signal of a few MeV but still gaining only by a square root)

where do we get if we do so ?

- zero background means you have only to be patient to accumulate the number of event you need

**expected
number of
 $\beta\beta 0\nu$ events**

$$S = \frac{M \cdot N_A \cdot a}{W} \cdot \ln(2) \cdot \frac{t}{T_{1/2}^{0\nu}} \cdot \varepsilon$$

detector mass isotopic abundance live time
molecular mass / \
 ββ0ν half-life

stretch isotopic abundance to 100% and assume the same for the efficiency , make the calculation for ${}^{76}\text{Ge}$

how patient should you be (to collect 3 say 3 events)

- ▷ for a half-life of 10^{28} y approximately 6 years
- ▷ being realistic with efficiency and duty cycle say 10 years
- ▷ in terms of neutrino masses, let me scale from GERDA result (5×10^{25} y, 0.15-0.33 eV). Gain 200 on half-life means improving on the mass limit as the square root so 10-25 meV at the lower bound of the inverted hierarchy prediction

something more worrisome

$$\mathcal{M} \equiv g_A^2 \mathcal{M}_{0\nu} = g_A^2 \left(M_{GT}^{(0\nu)} - \left(\frac{g_V}{g_A} \right)^2 M_F^{(0\nu)} + M_T^{(0\nu)} \right)$$

$$1/\tau = G(Q, Z) |M_{\text{nucl}}|^2 \langle M_{\beta\beta} \rangle^2$$

$$g_A = \begin{cases} g_{\text{nucleon}} &= 1.269 \\ g_{\text{quark}} &= 1 \\ g_{\text{phen.}} &= g_{\text{nucleon}} \cdot A^{-0.18} \end{cases}$$

} who knows ?
2νββ

$$g_A \rightarrow g_A \cdot (1 - \delta)$$

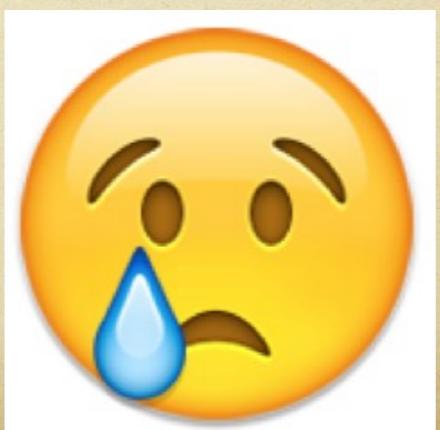
$$S \cdot (1 - \delta)^4$$

For instance, if we have a decrease by $\delta = 10$ (20)% of the axial coupling, lifetime would increase by a factor of $1/(1 - \delta)^4 = 1.5$ (2.5)

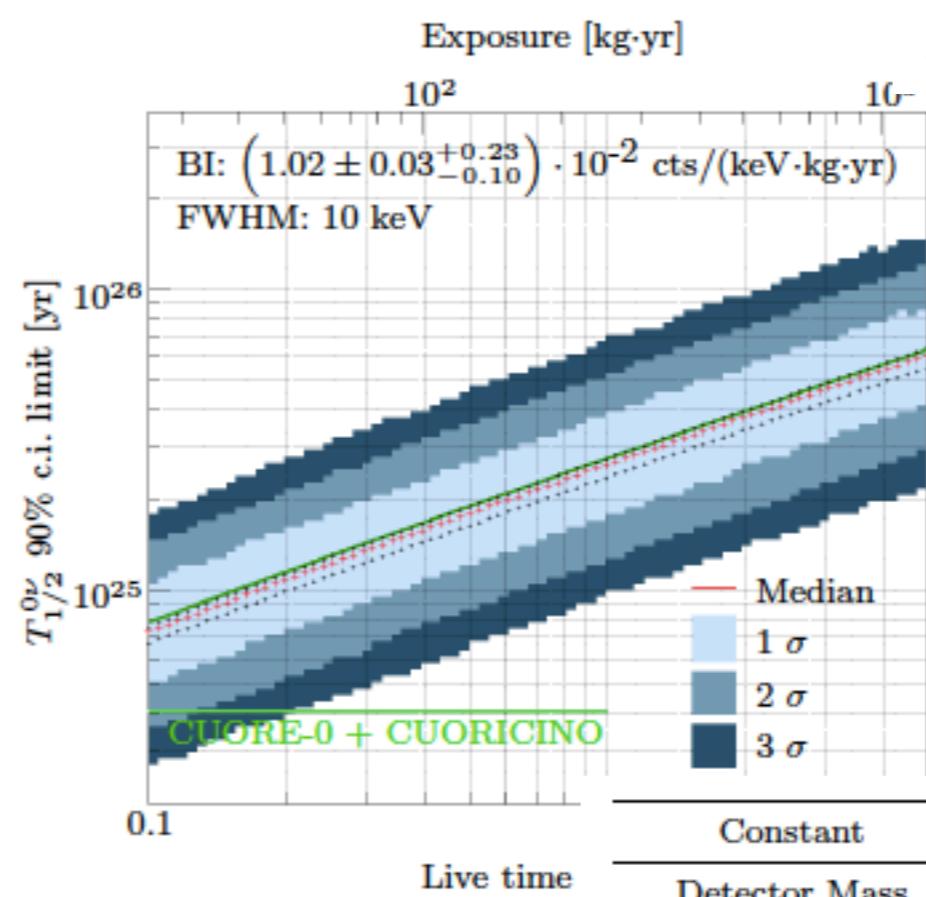
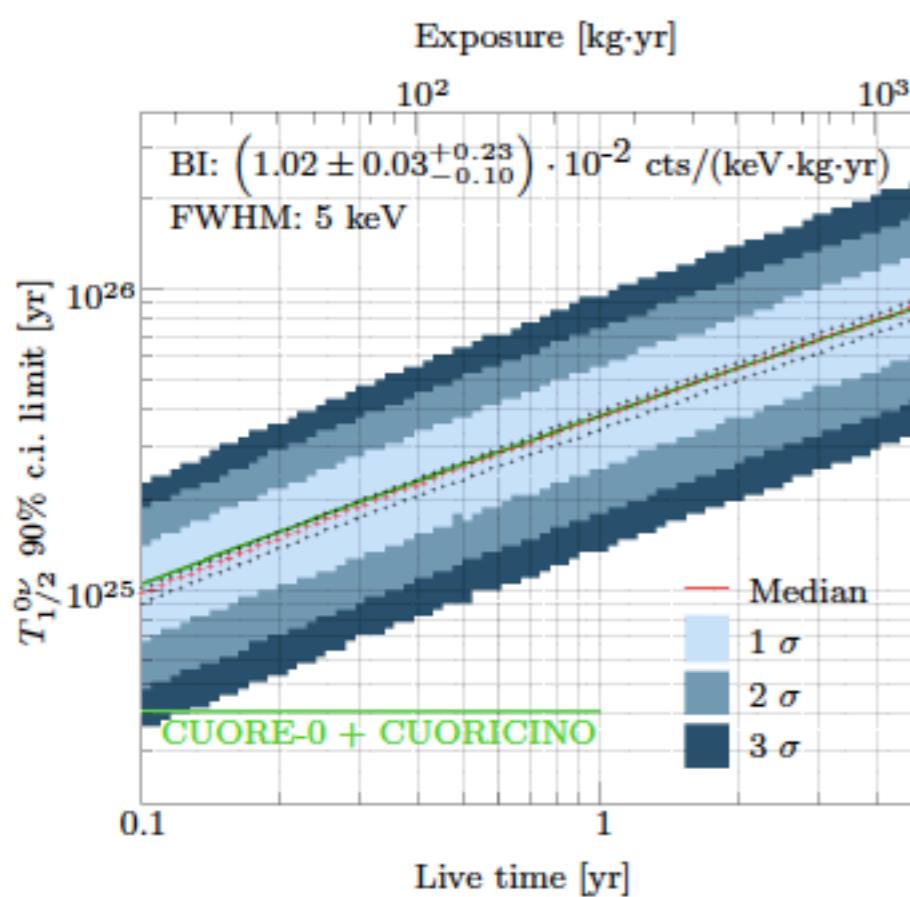
for ${}^{82}\text{Se}$
 $\delta = 0.55$
the ‘factor’
would be
11 !!!!!!

Conclusion

- ▷ worth to demonstrate the feasibility of zero background
- ▷ by achieving it, 10^{28} y would be a possible goal
- ▷ beyond that you only gain by mass increase, knowing however that the next challenge would be the direct ordering of neutrino masses which requires a jump of a factor 100 (100 tons ?)



CUORE SENSITIVITY



BI [cts/(keV·kg·yr)]
$(1.02 \pm 0.03(\text{stat})^{+0.23}_{-0.10}(\text{syst})) \cdot 10^{-2}$

Constant	Symbol	Value
Detector Mass	m_d	741.67 kg
Avogadro number	N_A	$6.022 \cdot 10^{23}$ mol ⁻¹
Molar mass	m_A	159.6 g/mol
Live Time	t_d	0.1–5 yr
Efficiency	ε_{tot}	81.3%
¹³⁰ Te abundance	f_{130}	0.34167
$0\nu\beta\beta$ Q-value	$Q_{\beta\beta}$	2527.518 keV
⁶⁰ Co peak position	μ_{Co}	(2505.692 ± 1.9) keV
Energy resolution	FWHM	5, 10 keV